

# Line Integrals

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# Differentials of Oriented Curves

Let  $C$  be an *oriented* curve (in  $\mathbb{R}^2$  or  $\mathbb{R}^3$ ), *parametrized* by

$$\mathbf{r}(t), \quad a \leq t \leq b \quad (\text{with } \mathbf{r}'(t) \neq \mathbf{0}).$$

We define the following *differentials* associated to  $C/\mathbf{r}$ .

- $ds = |\mathbf{r}'(t)| dt = |d\mathbf{r}| = \sqrt{dx^2 + dy^2 (+dz^2)}$
- $dx = x'(t) dt$ ,  $dy = y'(t) dt$  (and  $dz = z'(t) dt$ )
- $d\mathbf{r} = \mathbf{r}'(t) dt = \mathbf{i} dx + \mathbf{j} dy (+\mathbf{k} dz)$

These are used to integrate functions and vector fields along  $C$ .

## Example 1

## Example

Evaluate  $\int_C 4x^3 ds$ , where  $C$  is the portion of  $y = x^3 - 1$  in the fourth quadrant, oriented upward.

Since  $C$  is the graph of a function, we set  $x = t$  and  $y = t^3 - 1$ . We see from the graph that  $0 \leq t \leq 1$ . Now

$$dx = dt, \quad dy = 3t^2 dt \Rightarrow ds = \sqrt{1^2 + (3t^2)^2} dt = \sqrt{1 + 9t^4} dt.$$

Hence

$$\int_C 4x^3 ds = \int_0^1 4t^3 \sqrt{1 + 9t^4} dt = \left. \frac{2(1 + 9t^4)^{3/2}}{27} \right|_0^1 = \boxed{\frac{-2 + 20\sqrt{10}}{27}}.$$

## Example 2

## Example

Evaluate  $\int_C xy^2 dx$ ,  $\int_C xy^2 dy$  and  $\int_C xy^2 ds$  where  $C$  is the right half of  $x^2 + y^2 = 16$ , oriented counterclockwise.

Since  $C$  is part of a circle at the origin, we set  $x = 4 \cos t$  and  $y = 4 \sin t$ . We see from the graph that  $-\pi/2 \leq t \leq \pi/2$ . Now

$$dx = -4 \sin t dt, \quad dy = 4 \cos t dt \quad \Rightarrow \quad ds = 4 dt.$$

Hence

$$\int_C xy^2 ds = \int_{-\pi/2}^{\pi/2} 4^4 \cos t \sin^2 t dt = 4^4 \cdot \frac{\sin^3 t}{3} \Big|_{-\pi/2}^{\pi/2} = \boxed{\frac{2^9}{3}}.$$

## Example 2

(Continued)

Furthermore

$$\int_C xy^2 dx = -4^4 \int_{-\pi/2}^{\pi/2} \cos t \sin^3 t dt = -4^4 \cdot \frac{\sin^4 t}{4} \Big|_{-\pi/2}^{\pi/2} = \boxed{0}.$$

And finally

$$\begin{aligned} \int_C xy^2 dy &= 4^4 \int_{-\pi/2}^{\pi/2} \cos^2 t \sin^2 t dt = 4^4 \int_{-\pi/2}^{\pi/2} \left( \frac{\sin 2t}{2} \right)^2 dt \\ &= 4^3 \int_{-\pi/2}^{\pi/2} \frac{1 - \cos(4t)}{2} dt = 2^5 \cdot \left( t - \frac{\sin(4t)}{4} \right) \Big|_{-\pi/2}^{\pi/2} \\ &= \boxed{2^5 \pi}. \end{aligned}$$

## Example 3

## Example

Evaluate  $\int_C \langle x - y, xy \rangle \cdot d\mathbf{r}$ , where  $C$  is the portion of  $x^2 + y^2 = 4$  from  $(2, 0)$  to  $(0, -2)$ , counterclockwise.

Since  $C$  is part of a circle at the origin, we set  $x = 2 \cos t$  and  $y = 2 \sin t$ . We see from the graph that  $0 \leq t \leq 3\pi/2$ . Now

$$d\mathbf{r} = \langle -2 \sin t, 2 \cos t \rangle dt,$$

so that

$$\begin{aligned} \mathbf{F} \cdot d\mathbf{r} &= (2 \cos t - 2 \sin t)(-2 \sin t) + (4 \cos t \sin t)(2 \cos t) dt \\ &= -4 \sin t \cos t + 4 \sin^2 t + 8 \cos^2 t \sin t dt. \end{aligned}$$

## Example 3

(Continued)

Hence

$$\begin{aligned}\int_C \langle x - y, xy \rangle \cdot d\mathbf{r} &= \int_0^{3\pi/2} -4 \sin t \cos t + 4 \sin^2 t + 8 \cos^2 t \sin t \, dt \\ &= \boxed{\frac{2}{3} + 3\pi}.\end{aligned}$$

### Remarks:

- The first and last terms in the integral can be simultaneously antidifferentiated with the substitution  $u = \cos t$ .
- The middle term requires the half angle formula  $\sin^2 t = \frac{1 - \cos 2t}{2}$ .

# Remarks

## Parametrization Independence and Vanishing Tangent

- Any line integral  $\int_C$  depends on  $C$  only, *not* on the choice of  $\mathbf{r}(t)$  describing it.
- We can actually allow  $\mathbf{r}' = \mathbf{0}$  at finitely many points in  $[a, b]$ .
- One must make sure  $\mathbf{r}'$  doesn't change direction and retrace part of  $C$  at these points.
- The moral is that  $\mathbf{r}' \neq \mathbf{0}$  is not as important as  $\mathbf{r}$  being *one-to-one*.



# Remarks

## Some Basic Relationships

Let  $C$  be an oriented curve, starting at  $A$  and ending at  $B$ .

- $\int_C dx = x(B) - x(A)$  and likewise for  $dy$  and  $dz$ .
- Let  $-C$  denote  $C$  with the opposite orientation. Then

$$\int_{-C} f dx = - \int_C f dx, \quad \int_{-C} f ds = \int_C f ds,$$

$$\int_{-C} \mathbf{F} \cdot d\mathbf{r} = - \int_C \mathbf{F} \cdot d\mathbf{r}.$$

- These can be convenient when parametrizing in a particular direction is troublesome (e.g. a clockwise oriented circle).

# Piecewise Line Integrals

If  $C$  is made up of smooth pieces  $C_1, C_2, C_3, \dots$ , then

$$\int_C f \, dx = \int_{C_1} f \, dx + \int_{C_2} f \, dx + \int_{C_3} f \, dx \cdots,$$

and likewise for all other differentials.

## Example

*Evaluate  $\int_C xy \, dx + (x - y) \, dy$ , where  $C$  is the piecewise linear path from  $(0, 0)$  to  $(2, 0)$  to  $(3, 2)$ .*

## A Piecewise Example

$C$  consists of the line segments  $C_1$  from  $(0, 0)$  to  $(2, 0)$ , and  $C_2$  from  $(2, 0)$  to  $(3, 2)$ .

On  $C_1$ :  $y = 0$  and  $dy = 0$ . Hence  $\int_{C_1} xy \, dx + (x - y) \, dy = \boxed{0}$ .

On  $C_2$ :  $\mathbf{r}(t) = \langle 2, 0 \rangle + t\langle 1, 2 \rangle = \langle 2 + t, 2t \rangle$  with  $0 \leq t \leq 1$ .

Thus  $dx = dt$  and  $dy = 2 \, dt$ .

So

$$\begin{aligned} \int_{C_2} xy \, dx + (x - y) \, dy &= \int_0^1 (2 + t)(2t) + (2 + t - 2t)2 \, dt \\ &= \int_0^1 2t^2 + 2t + 4 \, dt = \frac{2}{3}t^3 + t^2 + 4t \Big|_0^1 = \boxed{\frac{17}{3}}. \end{aligned}$$

# Interpreting Differentials

Let  $f$  be a function and  $\mathbf{F}$  be a vector field. The line integrals

$$\int_C f \, ds, \quad \int_C f \, dx \text{ (or } dy \text{ or } dz), \quad \int_C \mathbf{F} \cdot d\mathbf{r}$$

can all be interpreted using Riemann sums.

- $ds$  represents an infinitesimal unit of *arclength* on  $C$ .
- $dx$  represents an infinitesimal change in  $x$  along  $C$ . Likewise with  $dy$  and  $dz$ .
- $d\mathbf{r}$  represents an infinitesimal displacement along  $C$ .

In certain situations these allow us to interpret the line integrals themselves.

# Arclength Integrals

If  $f = f(x, y)$ , then

$$\int_C f \, ds$$

is the (signed) area of the surface between  $C$  and the graph of  $f$  (a “fence”).

In particular, if  $f \equiv 1$ , then

$$\int_C ds = \text{arclength of } C.$$

## Example 4

### Example

Find the arclength of the portion of the helix  $\mathbf{r}(t) = \langle \cos t, \sin t, t \rangle$  between  $(1, 0, 0)$  and  $(1, 0, 2\pi)$ .

We have

$$d\mathbf{r} = \langle -\sin t, \cos t, 1 \rangle dt \Rightarrow ds = |d\mathbf{r}| = \sqrt{2} dt.$$

So the length is

$$\int_C ds = \int_0^{2\pi} \sqrt{2} dt = \boxed{2\pi\sqrt{2}}.$$

**Remark:** This can also be computed geometrically by “unrolling” the cylinder that the helix sits on.

# Line Integrals of Vector Fields

The differentials  $dx$ ,  $dy$ ,  $dz$  and  $d\mathbf{r}$  can be “interpreted” simultaneously.

If  $\mathbf{F} = P\mathbf{i} + Q\mathbf{j} + R\mathbf{k}$ , then on the one hand we have

$$\begin{aligned}\int_C \mathbf{F} \cdot d\mathbf{r} &= \int_C (P\mathbf{i} + Q\mathbf{j} + R\mathbf{k}) \cdot (\mathbf{i}dx + \mathbf{j}dy + \mathbf{k}dz) \\ &= \int_C P dx + Q dy + R dz.\end{aligned}$$

But we also have

$$\int_C \mathbf{F} \cdot d\mathbf{r} = \int_a^b \mathbf{F} \cdot \mathbf{r}'(t) dt = \int_a^b \frac{\mathbf{F} \cdot \mathbf{r}'(t)}{|\mathbf{r}'(t)|} |\mathbf{r}'(t)| dt = \int_C \text{proj}_{\mathbf{T}} \mathbf{F} ds,$$

where  $\mathbf{T}$  is the tangent vector to  $C$  at any point.

# Integrating $\text{proj}_{\mathbf{T}} \mathbf{F}$

What does the line integral

$$\int_C \mathbf{F} \cdot d\mathbf{r} = \int_C \text{proj}_{\mathbf{T}} \mathbf{F} \, ds$$

measure?

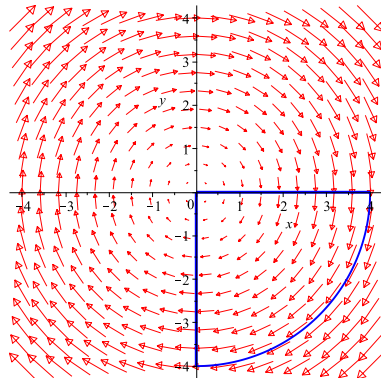
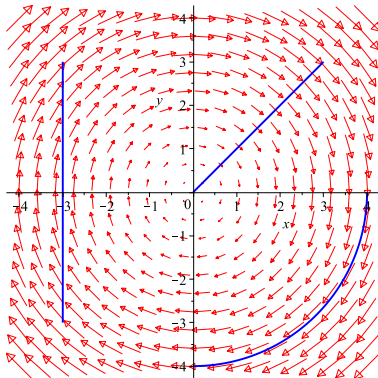
- Recall that  $\text{proj}_{\mathbf{T}} \mathbf{F}$  is:
  - \* positive when  $\mathbf{T}$  and  $\mathbf{F}$  tend to point in the same direction;
  - \* zero when  $\mathbf{T}$  and  $\mathbf{F}$  are orthogonal;
  - \* negative otherwise.
- The integral therefore measures the extent to which  $\mathbf{F}$  “points along”  $C$ .



# Example 5

## Example

For each curve  $C$  shown, choose an orientation and determine if  $\int_C \mathbf{F} \cdot d\mathbf{r}$  is positive, negative or zero.



# Work

If  $\mathbf{F}$  is a *force field*, we see that

$$\int_C \mathbf{F} \cdot d\mathbf{r} = \text{work done by } \mathbf{F} \text{ on a particle moving along } C.$$

If  $\mathbf{F}$  is the only force acting on the particle, then by Newton's second law,  $\mathbf{F} = m\mathbf{r}''$  and

$$\begin{aligned} \int_C \mathbf{F} \cdot d\mathbf{r} &= \int_a^b m\mathbf{r}''(t) \cdot \mathbf{r}'(t) dt = \frac{m}{2} \int_a^b \frac{d}{dt}(\mathbf{r}'(t) \cdot \mathbf{r}'(t)) dt \\ &= \frac{m}{2} (|\mathbf{r}'(b)|^2 - |\mathbf{r}'(a)|^2) = \frac{mv_f^2}{2} - \frac{mv_i^2}{2} = \Delta E_k, \end{aligned}$$

the change in kinetic energy from one end of  $C$  to the other.

# Integrating Conservative Fields

Suppose that  $\mathbf{F}$  is conservative with potential function  $f$ .

Then on any curve we have

$$\begin{aligned}\int_C \mathbf{F} \cdot d\mathbf{r} &= \int_C f_x dx + f_y dy + f_z dz \\ &= \int_a^b \frac{\partial f}{\partial x} \frac{dx}{dt} + \frac{\partial f}{\partial y} \frac{dy}{dt} + \frac{\partial f}{\partial z} \frac{dz}{dt} dt \\ &= \int_a^b \frac{df}{dt} dt = f(\mathbf{r}(b)) - f(\mathbf{r}(a)) \\ &= f(\text{end of } C) - f(\text{beg. of } C).\end{aligned}$$

This is the *Fundamental Theorem of Calculus for Line Integrals*.

# The Fundamental Theorem

## Theorem (Fundamental Theorem of Calculus for Line Integrals)

If  $\mathbf{F}$  is a conservative vector field with potential function  $f$ , then for any curve  $C$  from  $A$  to  $B$  we have

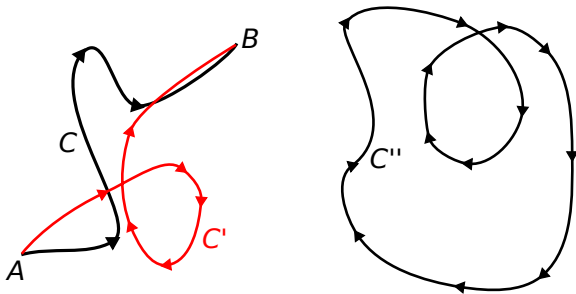
$$\int_C \mathbf{F} \cdot d\mathbf{r} = f(B) - f(A).$$

Consequently:

- $\int_C \mathbf{F} \cdot d\mathbf{r}$  is path independent; it depends only on the endpoints of  $C$ .
- $\int_C \mathbf{F} \cdot d\mathbf{r} = 0$  for any closed curve  $C$  (when  $A = B$ ).

# Path Independence

If  $\mathbf{F} = \nabla f$ , then  $\int_C \mathbf{F} \cdot d\mathbf{r} = \int_{C'} \mathbf{F} \cdot d\mathbf{r}$  and  $\int_{C''} \mathbf{F} \cdot d\mathbf{r} = 0$ .



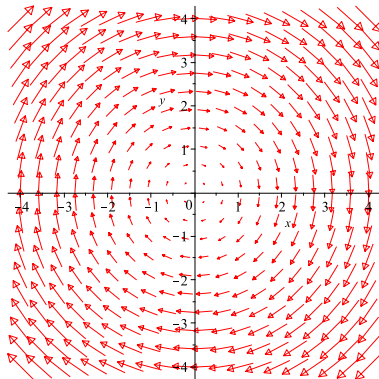
## Warning

*The FTC for Line Integrals only applies when  $\mathbf{F} = \nabla f$ . Not all  $\mathbf{F}$  have this property!*

# Example 6

## Example

*Explain why the vector field shown below is not conservative.*



## Example 7

### Example

Evaluate  $\int_C (x + 2xy) dx + x^2 dy$  where  $C$  is the curve consisting of the line segments from  $(0,0)$  to  $(2,1)$  to  $(4,3)$  to  $(5,0)$ .

Notice that if  $f = x^2y + \frac{x^2}{2}$ , then

$$\nabla f = \langle 2xy + x, x^2 \rangle,$$

which is the vector field under consideration.

So by the Fundamental Theorem

$$\int_C (x + 2xy) dx + x^2 dy = f(5,0) - f(0,0) = \boxed{\frac{25}{2}}.$$

## Example 8

## Example

Evaluate  $\int_C \mathbf{F} \cdot d\mathbf{r}$  where  $\mathbf{F} = x\mathbf{i} + (y + 2)\mathbf{j}$  and  $C$  is given by  $\mathbf{r}(t) = \langle t - \sin t, 1 - \cos t \rangle$ ,  $0 \leq t \leq 2\pi$ .

We could compute  $d\mathbf{r}$ , but instead we notice that

$$\nabla \underbrace{\left( \frac{x^2}{2} + \frac{y^2}{2} + 2y \right)}_f = x\mathbf{i} + (y + 2)\mathbf{j} = \mathbf{F}.$$

Hence

$$\int_C \mathbf{F} \cdot d\mathbf{r} = f(\mathbf{r}(2\pi)) - f(\mathbf{r}(0)) = f(2\pi, 0) - f(0, 0) = \boxed{2\pi^2}.$$



## Example 9

## Example

Evaluate  $\int_C \mathbf{F} \cdot d\mathbf{r}$  where  $\mathbf{F} = (y + z)\mathbf{i} + (x + z)\mathbf{j} + (x + y)\mathbf{k}$  and  $C$  is the line segment from  $(1, 0, 0)$  to  $(3, 4, 2)$ .

Not hard to parametrize a line segment, but again we have

$$\nabla \underbrace{(xy + xz + yz)}_f = (y + z)\mathbf{i} + (x + z)\mathbf{j} + (x + y)\mathbf{k} = \mathbf{F}.$$

So once more

$$\int_C \mathbf{F} \cdot d\mathbf{r} = f(3, 4, 2) - f(1, 0, 0) = \boxed{26}.$$

# Characterizing Conservative Fields

- Path independence of line integrals turns out to characterize conservative vector fields.
- If  $\int_C \mathbf{F} \cdot d\mathbf{r}$  is path independent in some domain, it is possible to construct a potential function by “integrating  $\mathbf{F}$  to  $X$ ”.
- This means the Fundamental Theorem is actually an “if and only if” result.

## Theorem

*A vector field  $\mathbf{F}$  with domain  $\Omega$  is conservative if and only if*

*$\int_C \mathbf{F} \cdot d\mathbf{r}$  is path independent for all  $C \subset \Omega$ .*

# Derivatives of Conservative Fields

Testing path independence for *all* curves can be difficult. How else might we identify a conservative field?

Notice that

$$\begin{aligned}\nabla \times \nabla f &= \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ f_x & f_y & f_z \end{vmatrix} \\ &= (f_{zy} - f_{yz})\mathbf{i} - (f_{zx} - f_{xz})\mathbf{j} + (f_{yx} - f_{xy})\mathbf{k} = \mathbf{0},\end{aligned}$$

by Clairaut's theorem.

## Moral

*If a field  $\mathbf{F}$  is conservative, then  $\nabla \times \mathbf{F} = \mathbf{0}$ .*

## Example 10

### A Nonconservative Field

The vector field  $\mathbf{F} = xy\mathbf{i} + yz\mathbf{j} + xz\mathbf{k}$  is *not* conservative since

$$\begin{aligned}\nabla \times \mathbf{F} &= \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ xy & yz & xz \end{vmatrix} \\ &= (-y)\mathbf{i} - (z)\mathbf{j} + (-x)\mathbf{k} \\ &= -(y\mathbf{i} + z\mathbf{j} + x\mathbf{k}) \\ &\neq \mathbf{0}.\end{aligned}$$

# Closed and Exact Forms

Consider a vector field  $\mathbf{F} = P\mathbf{i} + Q\mathbf{j} + R\mathbf{k}$ . The associated quantity

$$\omega = P dx + Q dy + R dz$$

is called a (*differential*) *1-form*.

Notice that we can write

$$\int_C \mathbf{F} \cdot d\mathbf{r} = \int_C P dx + Q dy + R dz = \int_C \omega.$$

$\omega$  is called *exact* when  $\mathbf{F} = \nabla f$ . We write  $\omega = df$

$\omega$  is called *closed* when  $\nabla \times \mathbf{F} = \mathbf{0}$ . We write  $d\omega = 0$ .

# Exact $\Rightarrow$ Closed, but...

We have seen that:

## Theorem

*Every exact form is closed.*

However:

## Warning

*Not every closed form is exact! Depending on the domain, there may exist  $\mathbf{F}$  that are not conservative, but for which  $\nabla \times \mathbf{F} = \mathbf{0}$ .*

That being said, as a consequence of Green's and Stokes' theorems:

## Theorem

*If  $\Omega$  is a simply connected domain (e.g.  $\Omega = \mathbb{R}^2$  or  $\Omega = \mathbb{R}^3$ ), then every closed form on  $\Omega$  is exact.*

# Example 11

## Example

Is  $\omega = y(z + 2) dx + (xz + 2x + 4) dy + (xy + 3) dz$  exact? If so, find a potential function.

Because  $\omega$  is defined on  $\mathbb{R}^3$ , and  $d\omega = 0$  (exercise),

$\omega$  is exact.

To find a potential we write

$$f_x = y(z + 2),$$

$$f_y = xz + 2x + 4,$$

$$f_z = xy + 3,$$

then repeatedly integrate and substitute.

# Example 11

## Continued

We have

$$f = \int f_x \partial x = \int y(z+2) \partial x = xy(z+2) + \underbrace{g(y,z)}_{\text{"constant"}}.$$

Now compare to the second and third equations:

$$\begin{aligned}xz + 2x + 4 = f_y = x(z+2) + g_y &\Rightarrow g_y = 4, \\xy + 3 = f_z = xy + g_z &\Rightarrow g_z = 3.\end{aligned}$$

Integrate  $g_y = 4$ :

$$g = \int g_y \partial y = \int 4 \partial y = 4y + \underbrace{h(z)}_{\text{"constant"}}.$$

And compare to  $g_z$ :

$$3 = g_z = h'(z) \Rightarrow h(z) = 3z + C.$$



## Example 11

Continued, continued

So the potentials of  $\omega$  are

$$f = xy(z+2) + g = xy(z+2) + 4y + h = \boxed{xy(z+2) + 4y + 3z + C}.$$

### Remarks:

- The constant  $C$  is arbitrary. Any choice will satisfy  $\omega = df$ .
- If  $\omega$  is known to be exact, the process above will *always* find  $f$ .
- In simple cases one can also find  $f$  by “guess and check.”

## Example 12

### A Non-Exact Closed Form

Consider the 1-form  $\omega = \frac{-y dx + x dy}{x^2 + y^2}$  on the domain  $\mathbb{R}^2 - \{\mathbf{0}\}$ , which is *not* simply connected.

Let  $C$  be the unit circle, traversed counterclockwise, which can be parametrized by

$$\mathbf{r}(t) = \cos t \mathbf{i} + \sin t \mathbf{j}, \quad 0 \leq t \leq 2\pi.$$

Thus  $dx = -\sin t dt$ ,  $dy = \cos t dt$  and  $x^2 + y^2 = 1$ , so that

$$\int_C \omega = \int_0^{2\pi} (-\sin t)(-\sin t) + (\cos t)(\cos t) dt = \int_0^{2\pi} 1 dt = 2\pi.$$

So  $\omega$  is *not* exact.

However

$$\begin{aligned}d\omega &= \frac{\partial}{\partial x} \left( \frac{x}{x^2 + y^2} \right) - \frac{\partial}{\partial y} \left( \frac{-y}{x^2 + y^2} \right) \\&= \frac{(x^2 + y^2) - x(2x)}{(x^2 + y^2)^2} + \frac{(x^2 + y^2) - y(2y)}{(x^2 + y^2)^2} \\&= \frac{2(x^2 + y^2) - 2x^2 - 2y^2}{(x^2 + y^2)^2} = 0,\end{aligned}$$

so that  $\omega$  is closed!

## (More) Work

If  $\mathbf{F} = \nabla f$  is a *conservative force field*, then  $E_p = -f$  is the *potential energy*. If  $C$  is any path from  $A$  to  $B$ ,

$$\int_C \mathbf{F} \cdot d\mathbf{r} = f(B) - f(A) = E_p(A) - E_p(B).$$

If a particle moves along  $C$  under the influence of  $\mathbf{F}$  only, by earlier work

$$E_k(B) - E_k(A) = \int_C \mathbf{F} \cdot d\mathbf{r} = E_p(A) - E_p(B),$$

or

$$E_p(A) + E_k(A) = E_p(B) + E_k(B).$$

That is, the *total energy*  $E = E_p + E_k$  is *conserved* by  $\mathbf{F}$ .

# Conservative Fields are ... Conservative!

This shows:

## Theorem

*A conservative force field obeys the Law of Conservation of Energy.*

Put another way:

## Moral

*The quantity "conserved" by a conservative vector field is the total energy.*